M09T.2

We're given the Hamiltonian

$$H = rac{p_1^2}{2m} + U(x_1) - fx_N + \sum_{i=2}^N rac{p_1^2}{2m} + U(x_n - x_{n_1}) \hspace{1.5cm} (1)$$

We need not worry about the momentum, in this problem. We calculate the partition function Z,

$$Z \propto \int e^{-eta H} = \int d^N x \exp \left[-eta U(x_1) + eta f x_N - eta \sum_{i=2}^N U(x_n - x_{n_1})
ight]$$

We can begin by first perfoming the integration of x_n , denoted by Z_n , which gives

$$egin{aligned} Z_n & \propto \int_{x_{n-1}}^{x_{n-1}+a} dx_n e^{eta u_0} e^{eta f x} + \int_{x_{n-1}+a}^{x_{n-1}+Q} dx_n e^{eta f x} \ & = rac{1}{eta f} \Big[e^{eta u_0} \left(e^{eta f(x_{n-1}+a)} - e^{eta f(x_{n-1})}
ight) + \left(e^{eta f(x_{n-1}+Q)} - e^{eta f(x_{n-1}+a)}
ight) \Big] \ & = rac{e^{eta f x_{n-1}}}{eta f} \left(e^{eta (fa + u_0)} + e^{eta f Q} - e^{eta u_0} - e^{eta f a}
ight). \end{aligned}$$

where we'll later take the limit $Q o \infty$.

From here, it is clear that after performing the integration over all coordinates (when $x_0=0$), we arrive at

$$Z \propto \left[rac{1}{eta f} \left(e^{eta (fa+u_0)} + e^{eta fQ} - e^{eta u_0} - e^{eta fa}
ight)
ight]^N = Z_0^N.$$

Interestingly, the partition function is the same as the single particle partition function raised to the power of N. Most of the work is now done. The average length of the system is given by

$$egin{align} \langle x_N
angle &= rac{\int d^N x \; x_N e^{-eta H}}{Z} \ &= rac{\int d^N x \; x_N \exp \left[-eta U(x_1) + eta f x_N - eta \sum_{i=2}^N U(x_n - x_{n_1})
ight]}{Z} \end{aligned}$$

However, note that

$$\int d^N x_N e^{-eta H} = rac{1}{eta} \int rac{\partial}{\partial f} \, e^{-eta H} = rac{1}{eta} rac{\partial}{\partial f} \int e^{-eta H} = rac{1}{eta} rac{\partial Z}{\partial f} \eqno(4)$$

So

$$egin{align*} \langle x_N
angle &= rac{N}{eta} rac{\left[rac{1}{eta f} \left(e^{eta (fa + u_0)} + e^{eta f Q} - e^{eta u_0} - e^{eta f a}
ight)
ight]^{N-1}}{\left[rac{1}{eta f} \left(e^{eta (fa + u_0)} + e^{eta f Q} - e^{eta u_0} - e^{eta f a}
ight)
ight]^{N}} rac{\partial Z_0}{\partial f} \ &= -rac{N}{eta f} + Na \, rac{e^{eta (fa + u_0)} - e^{eta f a}}{e^{eta (fa + u_0)} + e^{eta f Q} - e^{eta u_0} - e^{eta f a}} + rac{NQe^{eta f Q}}{e^{eta (fa + u_0)} + e^{eta f Q} - e^{eta u_0} - e^{eta f a}} \ &= -rac{N}{eta f} + Na + rac{e^{eta u_0}}{e^{eta (fa + u_0)} + e^{eta f Q} - e^{eta u_0} - e^{eta f a}} + rac{N(Q - a)e^{eta f Q}}{e^{eta (fa + u_0)} + e^{eta f Q} - e^{eta u_0} - e^{eta f a}} \ &= Na - rac{Nk_{
m B}T}{f} + rac{Na}{e^{eta f a} + e^{eta f (Q - u_0)} - e^{eta (fa - u_0)} - 1} \ &+ rac{N(Q - a)e^{eta f Q}}{e^{eta (fa + u_0)} + e^{eta f Q} - e^{eta u_0} - e^{eta f a}} \ \end{pmatrix}$$

If f < 0 , then taking the limit of $Q o \infty$ gives

$$\langle x_N
angle = Na - rac{Nk_{
m B}T}{f} + rac{Na}{e^{eta fa} - e^{eta (fa - u_0)} - 1} \hspace{1cm} (5)$$

otherwise the length diverges. As T o 0 , $\langle x_N
angle o Na$.



The idea is right.

In (5) when taking the T o 0 limit you forgot that f<0. Taking that into account you would get $\langle x_N
angle o 0$, which makes more sense than Na.

Also, you were supposed to consider the high- $\!T$ limit too.

Also there are several typos in your formulas which would be better to correct.