

J17M.1 (Coupled Pendulums)

(a)  \bar{e} coordinates for this problem are θ_1 & θ_2 .

\bar{e} positⁿ of mass 1 is: $\vec{r}_1 = \langle L \sin \theta_1, -L \cos \theta_1 \rangle$

\bar{e} positⁿ of mass 2 is: $\vec{r}_2 = \langle L \sin \theta_2, -L \cos \theta_2 \rangle$

\bar{e} Lagrangian is:

$$L = T - V$$

$$= \frac{1}{2} m_1 \dot{\vec{r}}_1^2 + \frac{1}{2} m_2 \dot{\vec{r}}_2^2 + mgL \cos \theta_1 + mgL \cos \theta_2 - \frac{1}{2} \kappa (\theta_2 - \theta_1)^2$$

$$= \frac{1}{2} m_1 L^2 \dot{\theta}_1^2 + \frac{1}{2} m_2 L^2 \dot{\theta}_2^2 + mgL \cos \theta_1 + mgL \cos \theta_2 - \frac{1}{2} \kappa (\theta_1^2 + \theta_2^2 - 2\theta_1 \theta_2)$$

$$\approx \frac{1}{2} m_1 L^2 \dot{\theta}_1^2 + \frac{1}{2} m_2 L^2 \dot{\theta}_2^2 - \frac{1}{2} mgL \theta_1^2 - \frac{1}{2} m_2 gL \theta_2^2 - \frac{1}{2} \kappa (\theta_1^2 + \theta_2^2 - 2\theta_1 \theta_2)$$

Let $\vec{\eta} = \begin{pmatrix} \theta_1 \\ \theta_2 \end{pmatrix} e^{-i\omega t}$

$$\Rightarrow L = \frac{1}{2} \vec{\eta}^T \underbrace{\begin{pmatrix} m_1 L^2 & 0 \\ 0 & m_2 L^2 \end{pmatrix}}_M \ddot{\vec{\eta}} - \frac{1}{2} \vec{\eta}^T \underbrace{\begin{pmatrix} m_1 gL + \kappa & -\kappa \\ -\kappa & m_2 gL + \kappa \end{pmatrix}}_K \vec{\eta}$$

To find \bar{e} normal frequencies, we solve $\det(K - \omega^2 M) = 0$ for ω^2 .

$$\begin{vmatrix} m_1 gL + \kappa - \omega^2 m_1 L^2 & -\kappa \\ -\kappa & m_2 gL + \kappa - \omega^2 m_2 L^2 \end{vmatrix} = 0.$$

$$(m_1 gL + \kappa - \omega^2 m_1 L^2)(m_2 gL + \kappa - \omega^2 m_2 L^2) - \kappa^2 = 0$$

$$m_1 m_2 g^2 L^2 + \kappa^2 + \omega^4 m_1 m_2 L^4 + m_1 gL\kappa + m_2 gL\kappa - m_1 m_2 gL^3 \omega^2 - m_1 m_2 gL^3 \omega^2 - m_1 \kappa L^2 \omega^2 - m_2 \kappa L^2 \omega^2 = 0$$

$$\omega^4 (m_1 m_2 L^4) + \omega^2 (-2m_1 m_2 gL^3 - m_1 \kappa L^2 - m_2 \kappa L^2) + (m_1 m_2 g^2 L^2 + m_1 gL\kappa + m_2 gL\kappa) = 0$$

$$\omega^2 = \frac{1}{2m_1 m_2 L^4} \left[2m_1 m_2 gL^3 + m_1 \kappa L^2 + m_2 \kappa L^2 \pm \left(4m_1^2 m_2^2 g^2 L^6 + m_1^2 \kappa^2 L^4 + m_2^2 \kappa^2 L^4 + 4m_1 m_2 g\kappa L^5 + 4m_1 m_2 g\kappa L^5 + 2m_1 m_2 \kappa^2 L^4 - 4m_1^2 m_2^2 g^2 L^6 - 4m_1^2 m_2 g\kappa L^5 - 4m_1 m_2^2 g\kappa L^5 \right)^{1/2} \right]$$

$$= \frac{1}{2m_1 m_2 L^4} \left(2m_1 m_2 gL^3 + m_1 \kappa L^2 + m_2 \kappa L^2 \pm \sqrt{m_1^2 \kappa^2 L^4 + 2m_1 m_2 \kappa^2 L^4 + m_2^2 \kappa^2 L^4} \right)$$

$$= \frac{1}{2m_1 m_2 L^4} \left[2m_1 m_2 gL^3 + m_1 \kappa L^2 + m_2 \kappa L^2 \pm (m_1 \kappa L^2 + m_2 \kappa L^2) \right]$$

$$\Rightarrow \omega^2 = \frac{g}{L} + \frac{\kappa}{L^2} \left(\frac{1}{m_1} + \frac{1}{m_2} \right) \quad \text{or} \quad \omega^2 = \frac{g}{L}$$

\bar{e} mode correspond^s to $\omega^2 = \frac{g}{L}$ must be \neq for which $\tau = 0$ (\bar{e} pendulums are in phase).

Thus, this normal mode has $\vec{\eta} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$

By orthogonality, \bar{e} other mode is $\vec{\eta} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$, where they are out of phase so \bar{e} for^s spr^s stores some energy.

(b) For either eigenvector, evolut^o is trivial: $\vec{\eta}_i(t) = \vec{\eta}_i(0) e^{i\omega_i t}$

Thus, for a given set of initial condit^os: $\theta_1(0)$ & $\theta_2(0)$

$$\begin{aligned} \vec{v}(0) = \begin{pmatrix} \theta_1(0) \\ \theta_2(0) \end{pmatrix} &= \frac{\theta_1(0)}{\sqrt{2}} (\vec{\eta}_1 + \vec{\eta}_2) + \frac{\theta_2(0)}{\sqrt{2}} (\vec{\eta}_1 - \vec{\eta}_2) \\ &= \frac{1}{\sqrt{2}} (\theta_1(0) + \theta_2(0)) \vec{\eta}_1 + \frac{1}{\sqrt{2}} (\theta_1(0) - \theta_2(0)) \vec{\eta}_2 \end{aligned}$$

$$\Rightarrow \vec{v}(t) = \begin{pmatrix} \theta_1(t) \\ \theta_2(t) \end{pmatrix} = \frac{1}{\sqrt{2}} (\theta_1(0) + \theta_2(0)) \vec{\eta}_1 e^{-i\omega_1 t} + \frac{1}{\sqrt{2}} (\theta_1(0) - \theta_2(0)) \vec{\eta}_2 e^{-i\omega_2 t}$$