

04 Jan 2022

2. Long Antenna Pattern

A thin, straight, conducting wire centered on the origin and oriented along the z -axis carries a current

$$I = \hat{z} I_0 \cos \omega t$$

everywhere along its length ℓ . This antenna will radiate electromagnetic waves with frequency ω and wavelength $\lambda = 2\pi c/\omega$. We will not assume that $\ell \ll \lambda$.

- (a) Because of the current, a time-dependent charge $q(t)$ will accumulate at the two ends of the wire. Give expressions for the charge and current *densities* $\rho(\vec{x}, t)$ and $\vec{j}(\vec{x}, t)$ on the wire. Show that the electric dipole moment of this charge distribution satisfies $p(t) = p_0 \sin(\omega t)$ and evaluate p_0
- (b) Use these source densities to construct the scalar and vector potentials everywhere outside the source region ($r \gg \ell$). Do not assume anything about the relative magnitudes of ℓ and λ . Do state the gauge you are using.
- (c) Compute the angular distribution of the energy flux radiated from this antenna. Show that it reduces to the standard electric dipole radiation pattern when $\lambda \gg \ell$. For general λ , show that the energy flux radiated perpendicular to the \hat{z} direction depends only on the maximum electric dipole moment p_0 (and agrees with the standard electric dipole radiation result).

$$a) \vec{I} = \hat{z} I_0 \cos(\omega t)$$

$$\vec{J} = \hat{z} I_0 \cos(\omega t) \delta(x) \delta(y)$$

$$I(t) = \frac{dq}{dt} \rightarrow q(t) = \frac{I_0}{\omega} \sin(\omega t)$$

$$\rho = q(t) \cdot \frac{1}{V} = \frac{I_0}{\omega} \sin(\omega t) \delta(x) \delta(y) \left(\delta(z - \frac{l}{2}) - \delta(z + \frac{l}{2}) \right)$$

$$P(t) = q(t) \cdot l = \frac{I_0 l}{\omega} \sin(\omega t) \rightarrow p_0 = \frac{I_0 l}{\omega}$$

$$\begin{aligned} b) V(\vec{r}, t) &= \frac{1}{4\pi\epsilon_0} \int d^3r' \frac{\rho(\vec{r}', t - |\vec{r} - \vec{r}'|/c)}{|\vec{r} - \vec{r}'|} \\ &= \frac{1}{4\pi\epsilon_0} \int d^3r' \frac{I_0}{\omega} \sin\left(\omega t - \frac{((x-x')^2 + (y-y')^2 + (z-z')^2)^{1/2}}{c/\omega}\right) \frac{\delta(x)\delta(y)(\delta(z-\frac{l}{2}) - \delta(z+\frac{l}{2}))}{((x-x')^2 + (y-y')^2 + (z-z')^2)^{1/2}} \\ &= \frac{1}{4\pi\epsilon_0} \int_{-\infty}^{\infty} dz \frac{I_0}{\omega} \sin\left(\omega t - \frac{(x^2 + y^2 + (z-z')^2)^{1/2}}{c/\omega}\right) \frac{\delta(z-\frac{l}{2})}{(x^2 + y^2 + (z-z')^2)^{1/2}} \\ &\quad - \frac{1}{4\pi\epsilon_0} \int_{-\infty}^{\infty} dz \frac{I_0}{\omega} \sin\left(\omega t - \frac{(x^2 + y^2 + (z-z')^2)^{1/2}}{c/\omega}\right) \frac{\delta(z+\frac{l}{2})}{(x^2 + y^2 + (z-z')^2)^{1/2}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{I_0}{\omega} \left[\sin\left(\omega t - \frac{(x^2 + y^2 + (z-\frac{l}{2})^2)^{1/2}}{c/\omega}\right) \frac{1}{(x^2 + y^2 + (z-\frac{l}{2})^2)^{1/2}} - \sin\left(\omega t - \frac{(x^2 + y^2 + (z+\frac{l}{2})^2)^{1/2}}{c/\omega}\right) \frac{1}{(x^2 + y^2 + (z+\frac{l}{2})^2)^{1/2}} \right] \\ &\quad \left| \begin{aligned} r_{\pm} &= (x^2 + y^2 + z^2 \pm 2zl - \frac{l^2}{4})^{1/2} = r \left(1 \pm \frac{z}{r} \frac{l}{r} - \frac{1}{4} \frac{l^2}{r^2} \right)^{1/2} \leftarrow x^2 + y^2 + z^2 \equiv r^2 \\ &= r \left(1 \pm \frac{l}{r} \cos\theta - \frac{1}{4} \frac{l^2}{r^2} \right)^{1/2} \approx r \left(1 \pm \cos\theta \frac{a}{r} \right) \leftarrow a = \frac{l}{2} \\ r_{\pm}^{-1} &= r^{-1} \left(1 \pm \cos\theta \frac{a}{r} \right)^{-1} \approx r^{-1} \left(1 \mp \frac{a}{r} \cos\theta \right) \approx \frac{1}{r} \leftarrow \frac{a}{r^2} \ll 1 \end{aligned} \right. \\ &\quad \downarrow K = \frac{\omega}{c} = \frac{1}{c/\omega} \\ &= \frac{1}{4\pi\epsilon_0} \frac{I_0}{\omega r} \left(\sin(\omega t - Kr_-) - \sin(\omega t - Kr_+) \right) \\ &= \frac{1}{4\pi\epsilon_0} \frac{I_0}{\omega r} \left[\sin(\omega t - kr + a \cos\theta) - \sin(\omega t - kr - a \cos\theta) \right] \\ &= \frac{1}{4\pi\epsilon_0} \frac{I_0}{\omega r} \left[\sin(\omega t - kr) \cos(ka \cos\theta) + \cos(\omega t - kr) \sin(ka \cos\theta) \right. \\ &\quad \left. - \sin(\omega t - kr) \cos(ka \cos\theta) + \cos(\omega t - kr) \sin(ka \cos\theta) \right] \\ &= + \frac{1}{4\pi\epsilon_0} \frac{I_0}{\omega r} 2 \cos(\omega t - kr) \sin(ka \cos\theta) \end{aligned}$$

$$V = + \frac{1}{2\pi\epsilon_0} \frac{I_0}{\omega r} \cos(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right)$$

$$\begin{aligned}\vec{A} &= \frac{\mu_0}{4\pi} \int d\vec{r}' \frac{\vec{j}(\vec{r}', t - |\vec{r} - \vec{r}'|/c)}{|\vec{r} - \vec{r}'|} \\ &= \frac{\mu_0}{4\pi} \int d\vec{r}' \hat{z} I_0 \cos\left(\omega t - \frac{((x-x')^2 + (y-y')^2 + (z-z')^2)^{1/2}}{c/\omega}\right) \frac{\delta(x)\delta(y)}{((x-x')^2 + (y-y')^2 + (z-z')^2)^{1/2}} \\ &= \frac{\mu_0}{4\pi} I_0 \hat{z} \int_{-a}^a \cos\left(\omega t - K(x^2 + y^2 + (z-z')^2)^{1/2}\right) \frac{1}{(x^2 + y^2 + (z-z')^2)^{1/2}} dz' \\ &\quad \left| \begin{aligned} r' &= (x^2 + y^2 + z^2 - 2zz' + z'^2)^{1/2} = r \left(1 - 2\frac{z'}{r} \cos\theta + \frac{z'^2}{r^2}\right)^{1/2} \approx r \left(1 - \frac{z'}{r} \cos\theta\right) \\ &= r - z' \cos\theta \end{aligned} \right. \\ &\quad \left| \begin{aligned} r'^{-1} &= \frac{1}{r} \left(1 - \frac{z'}{r} \cos\theta\right)^{-1} \approx \frac{1}{r} + \frac{z'}{r^2} \cos\theta \approx \frac{1}{r} \end{aligned} \right. \\ &= \frac{\mu_0}{4\pi} \frac{I_0}{r} \hat{z} \int_{-a}^a \cos(\omega t - Kr + Kz' \cos\theta) dz' \\ &= \frac{\mu_0}{4\pi} \frac{I_0}{r} \hat{z} \int_{-a}^a \cos(\omega t - Kr) \cos(Kz' \cos\theta) - \sin(\omega t - Kr) \sin(Kz' \cos\theta) dz' \quad \text{odd over interval} \\ &= \frac{\mu_0}{4\pi} \frac{I_0}{r} \cos(\omega t - Kr) \frac{1}{K \cos\theta} \left[\sin(Kz' \cos\theta) \right]_{-a}^a \\ &= \boxed{\frac{\mu_0 I_0}{2\pi r K \cos\theta} \cos(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) \hat{z}}$$

$$\text{Lorentz gauge: } \vec{\nabla} \cdot \vec{A} = \frac{\partial V}{\partial t}$$

$$c) \vec{E} = -\nabla V - \frac{\partial \vec{A}}{\partial t}, \vec{B} = \vec{\nabla} \times \vec{A}$$

$$\text{Spherical: } h_1 = 1, h_2 = r, h_3 = r \sin\theta$$

$$\nabla f = \frac{\partial f}{\partial r} \hat{r} + \frac{1}{r} \frac{\partial f}{\partial \theta} \hat{\theta} + \frac{1}{r \sin\theta} \frac{\partial f}{\partial \phi} \hat{\phi}, \vec{\nabla} \times \vec{A} = \frac{1}{r^2 \sin\theta} \begin{vmatrix} \hat{r} & r\hat{\theta} & r\sin\theta\hat{\phi} \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial \phi} \\ A_r & rA_\theta & r\sin\theta A_\phi \end{vmatrix}$$

$$\frac{\partial V}{\partial t} = 0, \vec{A} = \hat{z} = \hat{r} \cos\theta - \hat{\theta} \sin\theta \rightarrow A_\phi = 0 \quad (\text{also } \frac{\partial A_r}{\partial \phi} = \frac{\partial A_\theta}{\partial \phi} = 0)$$

$$\nabla V = \frac{\partial V}{\partial r} \hat{r} + \frac{1}{r} \frac{\partial V}{\partial \theta} \hat{\theta}, \vec{\nabla} \times \vec{A} = \frac{1}{r^2 \sin\theta} \begin{vmatrix} \hat{r} & r\hat{\theta} & r\sin\theta\hat{\phi} \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial \phi} \\ A_r & rA_\theta & 0 \end{vmatrix} = \frac{r \sin\theta}{r^2 \sin\theta} \left(\frac{\partial}{\partial r} (rA_\theta) - \frac{\partial}{\partial \theta} (A_r) \right) \hat{\phi}$$

$$\vec{A} \propto \frac{\hat{z}}{\cos\theta} = \hat{r} - \hat{\theta} \tan\theta$$

$$\vec{A} = \frac{\mu_0 I_0}{2\pi K} \left[\frac{1}{r} \cos(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) \right] (\hat{r} - \hat{\theta} \tan\theta)$$

$$A_r = \frac{\mu_0 I_0}{2\pi K} \left[\frac{1}{r} \cos(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) \right]$$

$$A_\theta = -\frac{\mu_0 I_0}{2\pi K} \left[\frac{1}{r} \cos(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) \right] \tan\theta$$

$$V = + \frac{1}{2\pi\epsilon_0} \frac{I_0}{\omega r} \cos(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right)$$

$$\begin{aligned} \frac{1}{r} \frac{\partial}{\partial r}(r A_\theta) &= -\frac{\mu_0 I_0 \tan\theta}{2\pi K r} \left[\cos(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) \right] \\ &= -\frac{\mu_0 I_0 \tan\theta}{2\pi K r} \left(-(-K) \sin(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) \right) \\ &= -\frac{\mu_0 I_0 \tan\theta}{2\pi V} \sin(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) \end{aligned}$$

$$\frac{1}{r} \frac{\partial}{\partial \theta}(A_r) = \frac{\mu_0 I_0 \tan\theta}{2\pi V^2} \frac{\partial}{\partial \theta} \left[\cos(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) \right]$$

$\approx 0 \leftarrow A_\theta$ goes as $\frac{1}{r}$ and will dominate for large r

$$\begin{aligned} \vec{\nabla} \times \vec{A} &= -\frac{\mu_0 I_0 \tan\theta}{2\pi V} \sin(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) \hat{\phi} \\ &= \vec{B} \end{aligned}$$

$$\begin{aligned} \frac{\partial V}{\partial r} &= -\frac{1}{2\pi\epsilon_0} \frac{I_0}{\omega r^2} \cos(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) + \frac{K}{2\pi\epsilon_0} \frac{I_0}{\omega r} \sin(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) \leftarrow \text{ignore } O(r^2) \\ &\approx \frac{K}{2\pi\epsilon_0} \frac{I_0}{\omega r} \sin(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) \end{aligned}$$

$$\frac{1}{r} \frac{\partial V}{\partial \theta} = \frac{1}{2\pi\epsilon_0} \frac{I_0}{\omega r^2} \cos(\omega t - Kr) \frac{\partial}{\partial \theta} \left[\sin\left(\frac{Kl}{2} \cos\theta\right) \right] \approx 0$$

$$\begin{aligned} -\nabla V &= -\frac{K}{2\pi\epsilon_0} \frac{I_0}{\omega r} \sin(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) \hat{r} \leftarrow \frac{K}{\omega} = \frac{1}{c} = \sqrt{\mu_0 \epsilon_0}, \frac{1}{\epsilon_0 c} = \sqrt{\frac{\mu_0}{\epsilon_0}} \\ &= -\frac{I_0}{2\pi V} \sqrt{\frac{\mu_0}{\epsilon_0}} \sin(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) \hat{r} \end{aligned}$$

$$\begin{aligned} -\frac{\partial \vec{A}}{\partial t} &= \frac{\mu_0 I_0}{2\pi K r} \sin\left(\frac{Kl}{2} \cos\theta\right) (\hat{r} - \hat{\theta} \tan\theta) \frac{\partial}{\partial t} [-\cos(\omega t - Kr)] \\ &= \frac{\mu_0 I_0}{2\pi r} \frac{\omega}{K} \sin(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) (\hat{r} - \hat{\theta} \tan\theta) \leftarrow \frac{\omega}{K} = c = \frac{1}{\sqrt{\epsilon_0 \mu_0}}, c \mu_0 = \sqrt{\frac{\mu_0}{\epsilon_0}} \\ &= \frac{I_0}{2\pi r} \sqrt{\frac{\mu_0}{\epsilon_0}} \sin(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) \hat{r} - \frac{I_0}{2\pi r} \sqrt{\frac{\mu_0}{\epsilon_0}} \sin(\omega t - Kr) \sin\left(\frac{Kl}{2} \cos\theta\right) \tan\theta \hat{\theta} \end{aligned}$$

$$\vec{E} = -\nabla V - \frac{\partial \vec{A}}{\partial t} = -\frac{I_0}{2\pi r} \sqrt{\frac{\mu_0}{\epsilon_0}} \sin(\omega t - kr) \sin\left(\frac{kl}{2} \cos\theta\right) \tan\theta \hat{\theta} \leftarrow \sqrt{\frac{\mu_0}{\epsilon_0}} = c\mu_0$$

$$= -\frac{I_0 \mu_0 c}{2\pi r} \sin(\omega t - kr) \sin\left(\frac{kl}{2} \cos\theta\right) \tan\theta \hat{\theta}$$

$$\vec{B} = \vec{\nabla} \times \vec{A} = -\frac{\mu_0 I_0 \tan\theta}{2\pi r} \sin(\omega t - kr) \sin\left(\frac{kl}{2} \cos\theta\right) \hat{\phi}$$

$$\vec{S} = \frac{1}{\mu_0} (\vec{E} \times \vec{B}) = \frac{\mu_0^2 c}{\mu_0} \left(\frac{I_0 \tan\theta}{2\pi r}\right)^2 \sin^2(\omega t - kr) \sin^2\left(\frac{kl}{2} \cos\theta\right) \hat{r}$$

$$= \mu_0 c \left(\frac{I_0 \tan\theta}{2\pi r}\right)^2 \sin^2(\omega t - kr) \sin^2\left(\frac{kl}{2} \cos\theta\right) \hat{r}$$

Energy flux distribution = $\frac{d}{d\Omega} \oint \vec{S} \cdot d\vec{a} = (\vec{S} \cdot \hat{r}) d\Omega = D(\theta)$

$$D(\theta) = \mu_0 c \left(\frac{I_0 \tan\theta}{2\pi r}\right)^2 \sin^2(\omega t - kr) \sin^2\left(\frac{kl}{2} \cos\theta\right) d\Omega$$

$$= \frac{\mu_0 c I_0^2}{(2\pi)^2 r^2} \sin^2(\omega t - kr) \left(\frac{kl}{2}\right)^2 \sin^2\theta \left(\frac{\sin^2(\frac{kl}{2} \cos\theta)}{\frac{kl}{2} \cos\theta}\right)^2 d\Omega$$

$$= \boxed{\frac{\mu_0 c I_0^2 k^2 l^2}{16 \pi^2 r^2} \sin^2(\omega t - kr) \sin^2\theta \operatorname{sinc}^2\left(\frac{kl}{2} \cos\theta\right) d\Omega}$$

$$\lambda \gg l, \frac{2\pi c}{\omega} \gg l \xrightarrow{k = \omega/c} kl \ll 2\pi, \frac{kl}{2} \ll \pi$$

$$\sin\left(\frac{kl}{2} \cos\theta\right) \approx 1$$

$$D(\theta) \approx \frac{\mu_0 c I_0^2 k^2 l^2}{16 \pi^2 r^2} \sin^2(\omega t - kr) \sin^2\theta d\Omega \leftarrow p_0 = \frac{I_0 l}{\omega} = \frac{I_0 l}{kc}$$

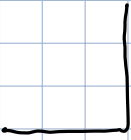
$$\approx \frac{\mu_0 c k^2}{16 \pi^2 r^2} p_0^2 (kc)^2 \sin^2(\omega t - kr) \sin^2\theta d\Omega$$

$$\approx \frac{\mu_0 p_0^2}{16 \pi^2 r^2} \left(c^3 \frac{\omega^4}{c^4}\right) \sin^2(\omega t - kr) \sin^2\theta d\Omega$$

$$\approx \frac{\mu_0 p_0^2 \omega^4}{16 \pi^2 r^2 c} \sin^2(\omega t - kr) \sin^2\theta d\Omega$$

$$\approx \boxed{\frac{\mu_0}{16 \pi^2 c} \left(p_0 \omega^2 \sin(\omega t - kr)\right)^2 \left(\frac{\sin\theta}{r}\right)^2 d\Omega}$$

This is the regular result of a dipole with $p(t) = p_0 \sin(\omega t - kr)$



There's another part but I can't be bothered, sorry :(