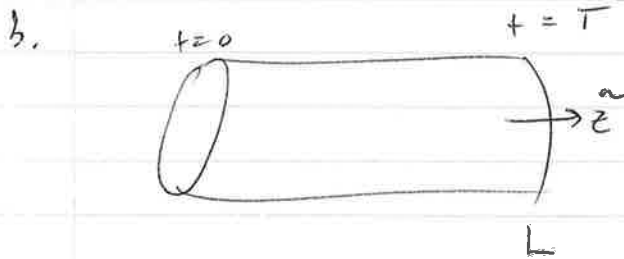


jlle 3

a.  $2\pi r B = \mu_0 \pi r^2 I$   
 $B = \frac{\mu_0 I}{2} r \hat{\phi}$



$$\vec{F} = q\vec{v} \times \vec{B} = \frac{qv\mu_0 I}{2} r (-\hat{r})$$

$$\ddot{r} = -\frac{qv\mu_0 I}{2m} r \rightarrow r = r_0 \cos at$$

$$\ddot{z} = -\frac{qv\mu_0 I}{2m} z \rightarrow z = -r_0 a \sin at$$

with  $a = \sqrt{\frac{qv\mu_0 I}{2m}}$

nonrelativistic, all exit tube at  $T = \frac{L}{v}$ .

at this time,  $r(T) = r_0 \cos aT$ ,  $z(T) = -r_0 a \sin aT$ .

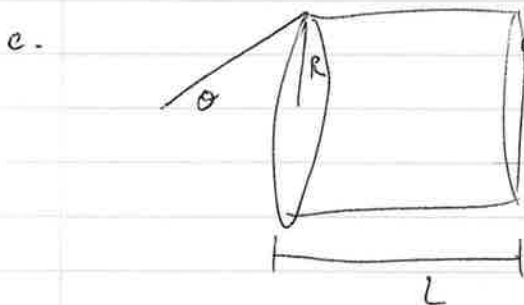
at some later time,  $r(t) = r(T) + \dot{r}(T)t + \dots$

$$0 = r_0 \cos aT - r_0 a \sin aT t + \dots$$

$$\rightarrow t' = \frac{\cos aT}{a \sin aT}$$

is when a particle reaches  $r=0$ , regardless of its initial radial position  $r_0$ .

the focal length is  $D = \frac{v}{\dot{z}} t' = \frac{v}{-a \cos aT} \frac{\cos aT}{a \sin aT} = \frac{v}{a^2 \sin aT}$   $a = \sqrt{\frac{qv\mu_0 I}{2m}}$



$$\theta \approx \tan \theta = \frac{R}{D} = \frac{R}{v} \sqrt{\frac{2m}{qv\mu_0 I}} \frac{\sin aT}{\cos aT}$$

$$\tan aT = \tan \left( L \sqrt{\frac{2\mu_0 I}{2m v}} \right) \approx \frac{L}{\sqrt{\dots}}$$