

## Prelims Solutions

### Problem J10T2

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#### 1

Consider a gas flowing in the  $x$  direction with  $v_x = v_x(z)$ . Viscosity  $\eta$  is defined by  $P_{zx} = -\eta \frac{\partial v_x}{\partial z}$  where  $P_{zx}$  is the change of the  $x$ -momentum in the  $z$  direction per unit time per unit area, which is a stress. Consider patch of area  $A$  in a  $z = \text{constant}$  plane and let  $n = \bar{n}, v = \bar{v}$ . The number of particles that flow through that patch in time  $\Delta t$  from below (or above) is  $\frac{1}{6}nvA\Delta t$  since  $1/3$  of particles will be going in the  $z$  direction and  $1/2$  of those will be moving in the plus (minus)  $z$  direction. A particle with mass  $m$  that came from below last had a collision roughly a mean free path  $l$  below the plane (i.e. at  $z - l$ ) and thus has  $x$  momentum  $mv_x(z - l)$ . A particle coming from above roughly has  $x$  momentum  $mv_x(z + l)$ . Thus,

$$P_{zx} = \frac{\frac{1}{6}nvA\Delta t m v_x(z - l) - \frac{1}{6}nvA\Delta t m v_x(z + l)}{A\Delta t} = -\frac{1}{3}nvlm \frac{\partial v_x}{\partial z}$$

Which gives  $\eta = \frac{1}{3}nvlm$  to first order in the Taylor expansion of  $v_x$ . Here I think the numerical factor in front shouldn't be trusted, but scaling should be correct.

#### 2

For hard spheres, the cross section doesn't change with temperature (it does in a plasma for ex.). Thus  $l = \frac{1}{n\sigma} \rightarrow \eta \approx \frac{mv}{\sigma} \approx \frac{(mk_B T)^{\frac{1}{2}}}{\sigma}$ . Thus the viscosity does not depend on pressure at constant  $T$ ! Interactions become non-negligible at higher pressures.

#### 3

Thermal wavelength is  $E = \frac{p^2}{2m} = \frac{h^2}{\lambda^2 m} \approx k_B T \rightarrow \frac{1}{\lambda} \approx \frac{\sqrt{mk_B T}}{h} \approx \frac{mv}{h}$  where  $mv^2 \approx k_B T$  from equipartition. So  $\eta \approx \frac{hnl}{\lambda}$ .

Using  $Z = \left(\frac{V}{\lambda^3}\right)^N / N!$  I get  $s = S/V = n(\ln(\frac{1}{n\lambda^3}) + \frac{5}{2})$

#### 4

$\frac{\eta}{s} = \frac{hl}{\lambda(\ln(\frac{1}{n\lambda^3}) + \frac{5}{2})} \approx h$  when all the scales become similar.