

1 January 2006, Quantum Mechanics, Problem 1

1.1 (a)

The first part was done by everyone, if you need to see a proof ask me:

$$\frac{d \langle \mathbf{r} \rangle}{dt} = \frac{\langle \mathbf{p} \rangle}{m} \quad (1)$$

$$\frac{d \langle \mathbf{p} \rangle}{dt} = - \langle \nabla V \rangle \quad (2)$$

1.2 (b)

Here there were a few solutions that I've seen. Hans suggested assuming that the field is small and solving to second order. I didn't think that was reasonable because they always say it in the statement when the field is supposed to be small. Then I heard some other solution, in which something was left as an infinite sum, or something like that. I didn't like that either. Finally, I came across a third solution, which is no more satisfactory in the sense that it requires some non-trivial assumption, but the final answer is very satisfying (to me). By steps:

Assumption 1. *The final state is a coherent state, and hence its wave function is:*

$$\Psi = B e^{iP_0 X} e^{-(X-X_0)^2}$$

where we are working with dimensionless units $X = \sqrt{\frac{m\omega}{\hbar}}x$ and $P = \frac{1}{\sqrt{m\hbar\omega}}p$.

1) Prove $\langle X \rangle = X_0$. This is done by writing down the integral and noting that it is of the form $\int X \text{Gaussian}(X - X_0)$, and thus is X_0 by parity arguments.

2) Prove $\langle P \rangle = P_0$. This is done by writing down $\langle \Psi | \hat{P} | \Psi \rangle$ and using the explicit form of the operator \hat{P} and the previous part.

3) Compute $\langle X \rangle$ and $\langle P \rangle$ by using part (a) of the problem.

4) Plug these values into X_0 and P_0 above, then check that it satisfies the initial condition that at time $t = 0$ the state should be the ground state of the harmonic oscillator with no electric field.

The final result is:

$$\Psi(x, t) = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{iqE\sin(\omega t)x/m\omega\hbar} e^{-\{x - \frac{qE}{m\omega^2}[1 - \cos(\omega t)]\}^2 m\omega/2\hbar} \quad (3)$$

As you may have guessed, the part I don't like is the assumption that the final state is a coherent state, because then you have to know this beforehand, and it is not trivial that one will.

The alternative way to solve this problem is like this:

$$|\Psi(x, t)\rangle = \sum_n \alpha_n e^{-iE_n t/\hbar} |n^+\rangle$$

where $+$ denotes "after turning on the electric field", not to be confused with the dagger that goes on the creation operator.

$$|\Psi(x, 0)\rangle = \sum_n \alpha_{n^+} |n^+\rangle = |0\rangle$$

$$\hat{T}(x_0) \sum_n \alpha_{n^+} |n^+\rangle = \hat{T}(x_0) |0\rangle = |0^+\rangle$$

where $x_0 = \frac{-qE}{m\omega^2}$, the amount by which the origin is shifted, and \hat{T} is the translation operator in position space.

$$\begin{aligned} \sum_n \alpha_{n^+} |n^+\rangle &= T^{\hat{-1}}(x_0) |0^+\rangle = \hat{T}(-x_0) |0^+\rangle \\ \langle m^+ | \sum_n \alpha_{n^+} |n^+\rangle &= \langle m^+ | \alpha_{m^+} |m^+\rangle = \alpha_{m^+} = \langle m^+ | \hat{T}(-x_0) |0^+\rangle \\ \alpha_{m^+} &= \langle m^+ | e^{ipx_0/\hbar} |0^+\rangle = \langle m^+ | e^{(-a^+ + a)x_0 \sqrt{m\omega/2\hbar}} |0^+\rangle \\ &= \langle m^+ | e^{-x_0 \sqrt{m\omega/2\hbar} a^+} e^{x_0 \sqrt{m\omega/2\hbar} a} e^{-m\omega x_0^2/4\hbar} |0^+\rangle \\ &= \langle m^+ | e^{-x_0 \sqrt{m\omega/2\hbar} a^+} |0^+\rangle e^{-m\omega x_0^2/4\hbar} = \langle m^+ | \sum_k \left(-x_0 \sqrt{\frac{m\omega}{2\hbar}} \right)^k \frac{(a^+)^k}{k!} |0^+\rangle \\ &= \langle m^+ | \left(-x_0 \sqrt{\frac{m\omega}{2\hbar}} \right)^m \frac{(a^+)^m}{m!} |0^+\rangle = \left(-x_0 \sqrt{\frac{m\omega}{2\hbar}} \right)^m \frac{1}{\sqrt{m!}} \\ |\Psi(x, t)\rangle &= \sum_n \left(-x_0 \sqrt{\frac{m\omega}{2\hbar}} \right)^n \frac{1}{\sqrt{n!}} e^{-i\left((n+\frac{1}{2})\hbar\omega - \frac{q^2 E^2}{2m\omega^2}\right)t/\hbar} |n^+\rangle \\ |\Psi(x, t)\rangle &= \sum_n \left(-x_0 \sqrt{\frac{m\omega}{2\hbar}} \right)^n \frac{1}{\sqrt{n!}} e^{-i\left((n+\frac{1}{2})\hbar\omega - \frac{q^2 E^2}{2m\omega^2}\right)t/\hbar} \frac{(a^+)^n}{\sqrt{n!}} |0^+\rangle \\ \Psi(x, t) &= e^{(-x_0 \sqrt{\frac{m\omega}{2\hbar}}) e^{-i\omega t} a^+} e^{-i\left(\frac{\hbar\omega}{2} - \frac{q^2 E^2}{2m\omega^2}\right)t/\hbar} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-m\omega\left(x + \frac{qE}{m\omega^2}\right)^2/2\hbar} \\ &\quad e^{ka^+} = e^{k(a^+ + a - a)} = e^{k(a^+ + a)} e^{-ka} e^{-\frac{k^2}{2}} \\ &\quad e^{ka^+} = e^{k(a^+ + a - a)} = e^{k\sqrt{\frac{2m\omega}{\hbar}}\left(x + \frac{qE}{m\omega^2}\right)} e^{-ka} e^{-\frac{k^2}{2}} \\ \Psi(x, t) &= e^{-x_0 \frac{m\omega}{\hbar} e^{-i\omega t} \left(x + \frac{qE}{m\omega^2}\right)} e^{-x_0^2 \frac{m\omega}{4\hbar} e^{-2i\omega t}} e^{-i\left(\frac{\hbar\omega}{2} - \frac{q^2 E^2}{2m\omega^2}\right)t/\hbar} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-m\omega\left(x + \frac{qE}{m\omega^2}\right)^2/2\hbar} \quad (4) \end{aligned}$$

This is a perfectly valid final answer, since it is in terms of x and t and known constants and functions. It would be nice, however, to check that it is the same as the previous answer. Also, the initial condition:

$$\begin{aligned} \Psi(x, 0) &= e^{\frac{qE}{\hbar\omega}\left(x + \frac{qE}{m\omega^2}\right)} e^{-\left(\frac{qE}{m\omega^2}\right)^2 \frac{m\omega}{4\hbar}} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-m\omega x^2/2\hbar} e^{-q^2 E^2/2\hbar m\omega^3} e^{-xqE/\hbar\omega} \\ \Psi(x, 0) &= e^{\frac{q^2 E^2}{4m\hbar\omega^3}} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-m\omega x^2/2\hbar} \end{aligned}$$

This just tells us that we made some mistake with a factor of 2, but we'll ignore that because the spirit of the solution is overall fine. If you actually check the two solutions against each other, they are not equal. I don't see the reason for this, so I was hoping to be enlightened by one of you.

Silviu provided yet another answer relying on Heisenberg picture operators. This solution gave a new answer, different from the other two. Clearly this problem leaves lots of things open to interpretation...