

J05Q.3

J05Q.3 - Heavy Particle Passing a Hydrogen Atom

According to time-dependent perturbation theory, the probability of a particle jumping from state k to k' is,

$$p_{k'k} = \frac{1}{\hbar^2} \left| \int_{-\infty}^{\infty} dt e^{\frac{i}{\hbar}(E_{k'} - E_k)t} \langle \psi_{k'} | H'(t) | \psi_k \rangle \right|^2 \quad (1)$$

Set the coordinate system with the hydrogen atom as the origin and trajectory of the heavy particle in x-z plane. The coordinate of the heavy particle can be denoted as $(vt, 0, D)$. Thus the distance between the heavy particle and the electron is,

$$s^2(t) = v^2 t^2 + D^2 - 2(vt \sin \theta) \cos \phi + 2D \cos \theta r + r^2 \quad (2)$$

As $D \gg a$,

$$\frac{1}{s(t)} \simeq \frac{1}{\sqrt{v^2 t^2 + D^2}} + (v^2 t^2 + D^2)^{-\frac{3}{2}} (vt \sin \theta \cos \phi + D \cos \theta) r \quad (3)$$

Thus the perturbation Hamiltonian is,

$$H' = -\frac{e^2}{4\pi\epsilon_0 s(t)} \quad (4)$$

Knowing that,

$$\psi_{10} = \frac{1}{\sqrt{\pi}} a^{-3/2} e^{-r/a}, \quad (5)$$

$$\psi_{210} = \frac{1}{4\sqrt{2}\pi} a^{-3/2} \frac{r}{a} e^{-r/2a} \cos\theta, \quad (6)$$

$$\psi_{21,\pm 1} = \frac{1}{8\sqrt{\pi}} a^{-3/2} \frac{r}{a} e^{-r/2a} \sin\theta e^{\pm i\phi}, \quad (7)$$

We can calculate,

$$\begin{aligned} & \langle \psi_{210} | H'(t) | \psi_{10} \rangle \\ &= -\frac{e^2}{4\pi\epsilon_0} \int_0^\infty r^2 dr \int_0^\pi \sin\theta d\theta \\ & \int_0^{2\pi} d\phi (v^2 t^2 + D^2)^{-\frac{3}{2}} D \cos\theta r \frac{1}{4\pi\sqrt{2}} a^{-4} r e^{-3r/2a} \cos\theta \\ &= -\frac{aDe^2}{12\sqrt{2}\pi\epsilon_0} (v^2 t^2 + D^2)^{-\frac{3}{2}} \int_0^\infty d\left(\frac{r}{a}\right) \left(\frac{r}{a}\right)^4 e^{-\frac{3r}{2a}} \\ &= -\frac{2^6 aDe^2}{3^5 \sqrt{2}\pi\epsilon_0} (v^2 t^2 + D^2)^{-\frac{3}{2}} \end{aligned}$$

$$\begin{aligned} & \langle \psi_{21,\pm 1} | H'(t) | \psi_{10} \rangle \\ &= -\frac{e^2}{4\pi\epsilon_0} \int_0^\infty r^2 dr \int_0^\pi \sin\theta d\theta \\ & \int_0^{2\pi} d\phi (v^2 t^2 + D^2)^{-\frac{3}{2}} vt \sin\theta \cos\phi r \frac{1}{8\pi} a^{-4} r e^{-3r/2a} \sin\theta e^{\pm i\phi} \\ &= -\frac{2^5 a e^2}{3^5 \pi \epsilon_0} vt (v^2 t^2 + D^2)^{-\frac{3}{2}} \end{aligned}$$

Then, as $\frac{D|E_0|}{\hbar v} \ll 1$

$$\begin{aligned} & \int_{-\infty}^\infty dt e^{\frac{i}{\hbar} \frac{3}{4} E_0 t} H'_{10 \rightarrow 210} \\ &= -\int_{-\infty}^\infty dt e^{\frac{i}{\hbar} \frac{3}{4} E_0 t} \frac{2^6 aDe^2}{3^5 \sqrt{2}\pi\epsilon_0} (v^2 t^2 + D^2)^{-\frac{3}{2}} \\ &\simeq -\int_{-\infty}^\infty dt \left(1 + \frac{i}{\hbar} \frac{3}{4} E_0 t\right) \frac{2^6 aDe^2}{3^5 \sqrt{2}\pi\epsilon_0} (v^2 t^2 + D^2)^{-\frac{3}{2}} \\ &= -\frac{2^6 a e^2}{3^5 \sqrt{2}\epsilon_0 D v} \end{aligned}$$

$$\begin{aligned}
& \int_{-\infty}^{\infty} dt e^{\frac{i}{\hbar} \frac{3}{4} E_0 t} H'_{10 \rightarrow 21, \pm 1} \\
&= - \int_{-\infty}^{\infty} dt e^{\frac{i}{\hbar} \frac{3}{4} E_0 t} \frac{2^5 a e^2}{3^5 \pi \epsilon_0} (vt) (v^2 t^2 + D^2)^{-\frac{3}{2}} \\
&\simeq - \int_{-\infty}^{\infty} dt \left(1 + \frac{i}{\hbar} \frac{3}{4} E_0 t\right) \frac{2^5 a e^2}{3^5 \pi \epsilon_0} (vt) (v^2 t^2 + D^2)^{-\frac{3}{2}} \\
&= - \frac{2^4 a e^2 i E_0}{3^5 \pi \epsilon_0 \hbar v^2}
\end{aligned}$$

Thus, the probability of the electron be in a $2p$ state is,

$$P_{10 \rightarrow 210} = \frac{2^{11} a^2 e^4}{\hbar^2 3^{10} \epsilon_0^2 D^2 v^2} \quad (8)$$

$$P_{10 \rightarrow 21, \pm 1} = \frac{2^8 a^2 e^4 E_0^2}{3^{10} \pi^2 \epsilon_0^2 \hbar^4 v^4} \quad (9)$$

One thought on "J05Q.3"



December 15, 2013 at 7:19 pm

All the steps are correct. If there are no computational mistakes, everything's correct so far. You're just missing one more step -- to find the total probability of $1s \rightarrow 2p$ transition.