

# 1 January 2005, Quantum Mechanics, Problem 1

As Hans pointed out to me, if you are asked to get the density matrix, you need the state of the system, because the density matrix is:

$$\rho = |\Psi\rangle\langle\Psi|$$

So we need to get the state of the system at an arbitrary time. This is relatively easy if you've done similar problems before. First, express the Hamiltonian in a more convenient form:

$$H = -J \left( \frac{S^2 - S_A^2 - S_B^2}{2} \right)$$

where  $S = S_A + S_B$ . The eigenstates of the system are eigenstates of definite value of total spin. These, can be expressed, in terms of the states of definite value of the z-component of each spin (where the + sign means an "up" state, and the order is always A first, then B):

$$\begin{aligned} |11\rangle &= |++\rangle \\ |10\rangle &= \frac{1}{\sqrt{2}}(|+-\rangle + |-+\rangle) \\ |1-1\rangle &= |--\rangle \\ |00\rangle &= \frac{1}{\sqrt{2}}(|+-\rangle - |-+\rangle) \end{aligned}$$

with respective energies:

$$\begin{aligned} E_{11} = E_{10} = E_{1-1} &= -\frac{J}{4} \\ E_{00} &= \frac{3J}{4} \end{aligned}$$

The initial state is  $|\Psi(0)\rangle = |+-\rangle = \frac{|10\rangle + |00\rangle}{\sqrt{2}}$ , and the state at time t is:

$$|\Psi(t)\rangle = Ae^{-iE_{00}t}|00\rangle + Be^{-iE_{11}t}|11\rangle + Ce^{-iE_{10}t}|10\rangle + De^{-iE_{1-1}t}|1-1\rangle$$

The initial condition determines the state as:

$$|\Psi(t)\rangle = \frac{1}{\sqrt{2}}[e^{-i3Jt/4}|00\rangle + e^{iJt/4}|10\rangle] = e^{-iJt/4}[\cos(Jt/2)|+-\rangle + i\sin(Jt/2)|-+\rangle]$$

The density matrix for the whole system is what I said in the beginning of the problem. But if we want the density matrix for A only, we have to use the formula, from Landau-Lifschitz:

$$\rho_A(x, x') = \int \Psi(q, x)\Psi^*(q, x')dq$$

where the interpretation is that  $x$  are the coordinates of A and  $q$  are the coordinates of B. In matrix notation, this can be written as:

$$\rho_A = \sum_{S_B} \langle S_B | S_A S_B \rangle \langle S'_A S_B | S_B \rangle$$

This is kind of hard to write in the language that we have, but it turns out that it is equal to:

$$\rho_A = \sum_{S_B} \langle S_B | \rho | S_B \rangle$$

When we write this we obtain:

$$\begin{aligned} \rho_A &= \cos^2(Jt/2) | \langle -_B | + - \rangle |^2 + \sin^2(Jt/2) | \langle +_B | - + \rangle |^2 \\ \rho_A &= \cos^2(Jt/2) | +_A \rangle \langle +_A | + \sin^2(Jt/2) | -_A \rangle \langle -_A | \end{aligned} \quad (1)$$

In matrix notation:

$$\rho_A = \begin{pmatrix} \cos^2(Jt/2) & 0 \\ 0 & \sin^2(Jt/2) \end{pmatrix} \quad (2)$$

This will describe a pure state whenever the average value of the z-component of the spin of A is either 1/2 or -1/2.

$$\langle S_z \rangle = Tr(S_z \rho_A) = \frac{\cos^2(Jt/2)}{2} - \frac{\sin^2(Jt/2)}{2}$$

Thus, we want the argument of the sine and cosine to be an integer multiple of  $\pi/2$ :

$$t = \frac{n\pi}{J} \quad (3)$$