

J02Q.3

Solution to J02Q.3 — Interacting Particles on a Line

Part a

First change variables to the center of mass frame

$$\begin{aligned}\bar{x} &= \frac{x_1 + x_2}{2} \\ x &= x_1 - x_2\end{aligned}\tag{1}$$

Making the Hamiltonian

$$H = \frac{p^2 + 4\bar{p}^2}{m} + (2\hbar^2 - S_T^2)U_0(x)\tag{2}$$

In the center of mass frame $\bar{p} = 0$ so we consider just the Hamiltonian

$$H = \frac{p^2}{m} + (2\hbar^2 - S_T^2)U_0(x)\tag{3}$$

We can separate the wavefunction into position and spin parts

$$\psi = \psi_x \psi_{S_A} \psi_{S_B}\tag{4}$$

Since the Hamiltonian is only concerned with S_T^2 , the eigenfunctions for the spin wavefunction will be the standard triplet-singlet states for two spin-1/2 particles. Written in the S,m basis they are:

$$\psi_{S_A} \psi_{S_B} = |1, 1\rangle, |1, 1\rangle, |1, -1\rangle, |0, 0\rangle\tag{5}$$

The Hamiltonian now splits into two cases: one for $S = 1$, where $S_T^2 = \hbar^2(S)(S+1) = 2\hbar^2$ and another for $S = 0$, where $S_T^2 = 0$.

$$H(S=1) = \frac{p^2}{m} \quad (6)$$

$$H(S=0) = \frac{p^2}{m} + 2\hbar^2 U_0(x)$$

The solution to the S=1 Hamiltonian trivially that of a free particle

$$\psi_x(S=1) = C e^{ik_1 x} \quad (7)$$

$$E(S=1) = \frac{\hbar^2 k_1^2}{m} \quad (8)$$

Where k_1 can take any real value. Thus when S=1 we do not get a bound state.

For the S=0 Hamiltonian we solve the Schrodinger equation:

$$\begin{aligned} H(S=0)\psi_x &= \frac{p^2}{m}\psi_x + 2\hbar^2 U_0(x)\psi_x \\ E\psi_x &= -\frac{\hbar^2}{m} \frac{\partial^2 \psi_x}{\partial x^2} + 2\hbar^2 \left(-\frac{\pi^2}{4ma^2}\right)\psi_x \\ \frac{\partial^2 \psi_x}{\partial x^2} &= -\left(\frac{mE}{\hbar^2} + \frac{\pi^2}{2a^2}\right)\psi_x \end{aligned}$$

The solution to this, using the condition imposed by the potential that $\psi_x = 0$ at $x = a$, is simply

$$\psi_x(S=0) = D \sin(\omega_0(x-a)) \quad (9)$$

Where

$$\omega_0^2 = \left(\frac{mE}{\hbar^2} + \frac{\pi^2}{2a^2}\right) \quad (10)$$

The potential also imposes the constraint $\psi_x = 0$ at $x = -a$. Plugging this into our solution gives the relation $\omega_0 = \frac{n\pi}{2a}$. We can now plug this into equation (7) and solve for the energy.

$$E_n(S=0) = \frac{\hbar^2 \pi^2 (n^2 - 2)}{4ma^2} \quad (11)$$

The ground state energy ($n=1$) is:

$$E_{n=1} = \frac{\hbar^2 \pi^2}{4ma^2} \quad (12)$$

Finally, normalizing the wave function gives $D = \sqrt{\frac{2}{\pi a}}$, leaving us with

$$\psi_x(S=0) = \sqrt{\frac{2}{\pi a}} \sin\left(\frac{\pi(x-a)}{2a}\right) \quad (13)$$

Part b

For a magnetic field $B_z = B_0 \cos k(x - ct)$ the perturbation Hamiltonian is

$$\begin{aligned} H' &= \mu \mathbf{B}(x_1) \cdot \mathbf{S}_A + \mu \mathbf{B}(x_2) \cdot \mathbf{S}_B \\ H' &= \mu B_z(x_1) S_{A_z} + \mu B_z(x_2) S_{B_z} \\ H' &= \mu B_z \left(\frac{x_1+x_2}{2} + \frac{x_1-x_2}{2} \right) S_{A_z} + \mu B_z \left(\frac{x_1+x_2}{2} - \frac{x_1-x_2}{2} \right) S_{B_z} \end{aligned}$$

We are given $ka \ll 1$. Since $|x_1 - x_2|$ is less than a , this implies that $k \frac{x_1-x_2}{2} \ll 1$, so we can Taylor expand the above expression to get

$$H' = \mu B_z \left(\frac{x_1 + x_2}{2} \right) S_{T_z} + \frac{x_1 - x_2}{2} \mu B'_z \left(\frac{x_1 + x_2}{2} \right) (S_{A_z} - S_{B_z}) \quad (14)$$

From part (a), the ground state had $S=0$. The spin wave function corresponding to this is

$$\psi_{S_A} \psi_{S_B} = |0,0\rangle = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \quad (15)$$

Apply the operators in our perturbation to this eigenstate:

$$S_{T_z} |0,0\rangle = 0 \quad (16)$$

$$(S_{A_z} - S_{B_z}) \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle) = |1,0\rangle \quad (17)$$

So we can drop the first term in the Hamiltonian, and we know that the bound state can only transition into the $|1,0\rangle$ state. To compute this probability we use time-dependent perturbation theory.

$$c_1(t) = -\frac{i}{\hbar} \int_0^t dt' e^{i \frac{E(S=1) - E_0}{\hbar} t'} \langle \psi_x(S=1) | \frac{x_1 - x_2}{2} \mu B'_z \left(\frac{x_1 + x_2}{2} \right) | \psi_x(S=0) \rangle \quad (18)$$

One thought on “J02Q.3”



December 15, 2013 at 10:32 pm

You're on the right track, try to finish your solution.

There are some typos and small mistakes in what you've written so far.

In (2) you have the wrong factor before \bar{p}^2 (not important for the solution though). In (7) forgot the i factor. In (10) and in the equation before it you've lost the factor of **2**, so that actually $\omega_0^2 = \left(\frac{mE}{\hbar^2} + \frac{\pi^2}{2a^2} \right)$ and the ground state energy is negative. Everything else seems to be fine.
