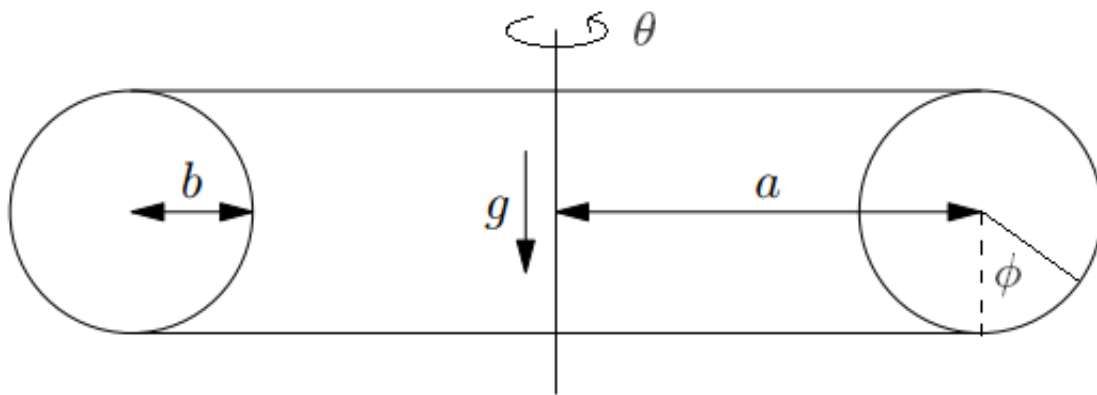


# J01M.1

## Problem

Find the frequency of small oscillations about uniform circular motion of a point mass that is constrained to move on the surface of a torus (donut) of major radius  $a$  and minor radius  $b$  whose axis is vertical.



Consider the circular cross section of the right side of the torus. Let  $\phi$  be the angle the point mass makes with the vertical. Let  $r$ ,  $z$ , and  $\theta$  be the standard cylindrical coordinates of the point mass relative to the axis of the torus.

$$\begin{aligned} z &= -b \cos \phi \\ \dot{z} &= b \sin \phi \dot{\phi} \\ r &= a + b \sin \phi \\ \dot{r} &= b \cos \phi \dot{\phi} \end{aligned}$$

Now we form the Lagrangian:

$$L = \frac{1}{2} m \dot{z}^2 + \frac{1}{2} m \dot{r}^2 + \frac{1}{2} m r^2 \dot{\theta}^2 - mgz$$

$$L = \frac{1}{2} b^2 \dot{\phi}^2 + \frac{1}{2} m(a + b \sin \phi)^2 \dot{\theta}^2 + gb \cos \phi$$

The Euler-Lagrange Equations for  $\theta$  and  $\phi$ .

$$(1) m(a + b \sin \phi) \dot{\theta} = \text{const} = l$$

$$(2) b\ddot{\phi} = (a + b \sin \phi) \dot{\theta}^2 \cos \phi - g \sin \phi$$

Equilibrium occurs when  $\ddot{\phi} = 0$  and  $\dot{\theta} = \Omega$ . Then (2) gives the equilibrium angle  $\phi_0$ :

$$(3) g \tan \phi_0 = (a + b \sin \phi_0) \Omega^2$$

And (1) gives us a value for  $l$

$$(4) l = m(a + b \sin \phi_0) \Omega$$

Eliminate  $\dot{\theta}$  from (2) using (1)

$$b\ddot{\phi} = \frac{l^2}{m^2(a+b \sin \phi)} \cos \phi - g \sin \phi$$

Introduce a small displacement  $\phi_0 \rightarrow \phi_0 + \epsilon$ .

$$b\ddot{\epsilon} = \frac{l^2}{m^2(a+b \sin(\phi_0+\epsilon))} \cos(\phi_0 + \epsilon) - g \sin(\phi_0 + \epsilon)$$

Expanding the sines and cosines, then using the small angle approximation leads to:

$$b\ddot{\epsilon} = \frac{l^2}{m^2(a+b(\sin \phi_0 + \epsilon \cos \phi_0))} (\cos \phi_0 - \epsilon \sin \phi_0) - g(\sin \phi_0 + \epsilon \cos \phi_0)$$

Taylor expanding the denominator

$$b\ddot{\epsilon} = \frac{l^2}{m^2(a+b \sin \phi_0)} \left( \cos \phi_0 - \epsilon \sin \phi_0 - \frac{b\epsilon \cos^2 \phi_0}{a+b \sin \phi_0} \right) - g(\sin \phi_0 + \epsilon \cos \phi_0)$$

Observe that the terms independent of  $\epsilon$  are the equilibrium condition and sum to 0 giving us:

$$b\ddot{\epsilon} = - \left( \frac{l^2}{m^2(a+b \sin \phi_0)} \left( \sin \phi_0 + \frac{b\epsilon \cos^2 \phi_0}{a+b \sin \phi_0} \right) + g \cos \phi_0 \right) \epsilon$$

Use (4) to eliminate  $l$  and then (3) to eliminate  $\Omega$

$$\ddot{\epsilon} = - \frac{g}{b} \left( \tan \phi_0 \left( \sin \phi_0 + \frac{b\epsilon \cos^2 \phi_0}{a+b \sin \phi_0} \right) + \cos \phi_0 \right) \epsilon$$

Therefore the frequency of small oscillations is:

$$\omega^2 = \frac{g}{b} \left( \tan \phi_0 \left( \sin \phi_0 + \frac{b\epsilon \cos^2 \phi_0}{a+b \sin \phi_0} \right) + \cos \phi_0 \right)$$

## 3 thoughts on "J01M.1"



OK.

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Approach is correct. However something should be fixed.

Note that a particle on a torus has two degrees of freedom. But your Lagrangian has only one degree of freedom  $\phi$ , because you substituted  $\Omega$  for  $\dot{\theta}$  (which is not legitimate).  $\dot{\theta}$  is not a conserved quantity of the problem (it changes as  $\phi$  changes). You indeed can eliminate  $\dot{\theta}$ , but not in such a way.

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I have updated my solution to use angular momentum to eliminate  $\dot{\theta}$

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