

Jan 2001 #2 (EM)

Relativistic cyclotron motion:

$$\frac{d\vec{p}}{dt} = \frac{q\vec{v} \times \vec{B}}{c} \quad \frac{d(m\gamma\vec{v})}{dt} = \frac{q\vec{v} \times \vec{B}}{c} \quad \text{For pure circular motion at constant speed, } \gamma = \text{constant}$$

\Rightarrow same solution as non-relativistic cyclotron motion with $m \rightarrow m\gamma$

$$\Omega = \frac{qB}{m\gamma c} \quad |\Omega| = \frac{eB}{m\gamma c}$$

$$R = \frac{v}{|\Omega|} = \frac{c}{e} \frac{m\gamma v}{B} = \frac{c}{e} \frac{p}{B} \quad p = |\vec{p}| \quad \text{magnitude of } \vec{p}$$

The problem is asking what the condition on B is to make R a constant

$$\Rightarrow \frac{dR}{dt} = 0 \quad \frac{1}{B} \frac{dp}{dt} - \frac{p}{B^2} \frac{dB}{dt} = 0 \quad \frac{B}{p} \frac{dp}{dt} = \frac{dB}{dt}$$

$\frac{dp}{dt}$? \vec{p} changes magnitude because of an inductive electric field, not the magnetic field directly:

$$\frac{dp}{dt} = -eE_{\phi}$$

E_{ϕ} is determined by Faraday's law:

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t} \Rightarrow \int \vec{E} \cdot d\vec{l} = -\frac{1}{c} \frac{d}{dt} \int \vec{B} \cdot d\vec{A}$$

$$2\pi R E = -\frac{\pi R^2}{c} \frac{dB_{z-\text{ave}}}{dt}$$

$$E = -\frac{R}{2c} \frac{dB_{z-\text{ave}}}{dt}$$

$$\Rightarrow \frac{dp}{dt} = \frac{eR}{2c} \frac{dB_{z-\text{ave}}}{dt}$$

$$\Rightarrow \frac{eR}{2c} \frac{B}{p} \frac{dB_{z-\text{ave}}}{dt} = \frac{dB}{dt} \quad R = \frac{c}{e} \frac{p}{B}$$

$$\Rightarrow \frac{c}{e} \frac{p}{B} \cdot \frac{e}{2c} \frac{B}{p} \frac{dB_{z-\text{ave}}}{dt} = \frac{dB}{dt}$$

$$\Rightarrow \frac{1}{2} \frac{dB_{z-\text{ave}}}{dt} = \frac{dB}{dt}$$

$$\frac{d}{dt} \left(B - \frac{B_{z-\text{ave}}}{2} \right) = 0 \quad B = \frac{B_{z-\text{ave}}}{2} + \text{const.}$$

maximum energy: power in = power radiated

$$P_{\text{rad}} = \frac{2}{3} \frac{e^2}{m^2 c^3} \left[\left(\frac{d\vec{p}}{dt} \right)^2 - \frac{1}{c^2} \left(\frac{dE}{dt} \right)^2 \right]$$

For relativistic circular motion, $\left| \frac{d\vec{p}}{dt} \right| = \gamma \Omega |\vec{p}| \gg \frac{1}{c} \frac{dE}{dt}$ $|\vec{p}| = m\gamma v$

$$P_{\text{rad}} = \frac{2}{3} \frac{e^2}{m^2 c^3} \cdot \gamma^2 \Omega^2 \cdot m^2 \gamma^2 v^2 = \frac{2}{3} \frac{e^2 \Omega^2 v^2 \gamma^4}{c^3}$$

$$\Omega = \frac{v}{R} \quad \text{or} \quad \Omega = \frac{eB}{m\gamma c}$$

ultra-relativistic: $v \approx c$; $\Rightarrow R \approx \frac{E}{eB}$ ($E \gg mc^2$) (energy $E = \gamma mc^2$)

$$\Rightarrow P_{\text{rad}} = \frac{2}{3} \frac{e^2}{m^4 c^7} \frac{E^4}{R^2} = \frac{2}{3} \frac{e^2}{m^4 c^7} e^2 B^2 E^2 \quad (\text{equivalent forms})$$

$$P_{\text{in}} = -e\vec{E} \cdot \vec{v} = \frac{eR}{2} \frac{dB_{z-\text{ave}}}{dt}$$

maximum E when $P_{\text{rad}} = P_{\text{in}}$

$$\frac{2}{3} \frac{e^2}{m^4 c^7} \frac{E^4}{R^2} = \frac{eB}{2} \frac{dB_{z-\text{ave}}}{dt}$$

$$E^4 = \frac{3}{4} \frac{m^4 c^7}{e} R^2 \frac{dB_{z-\text{ave}}}{dt}$$

$$\text{or } E = \frac{3}{4} \frac{m^4 c^7}{e^4} \frac{1}{B^2} \frac{dB_{z-\text{ave}}}{dt}$$